

Black Hole Entropy in the presence of Chern-Simons terms

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1. Introduction

2. Black Hole Spacetime

3. Black Hole Spacetime in Higher Derivative Gravity

4. Derivation of the First law

5. Summary

Black Holes Thermodynamics

- Classical General Relativity leads to

$$\frac{\kappa}{2\pi} \frac{\delta A}{4G_N} = \delta m - \Omega \delta J$$

- Semiclassical analysis identifies

$$T_H = \frac{\kappa}{2\pi}.$$

- Very natural to identify

$$S = \frac{A}{4G_N}$$

and look for statistical explanation.

- Extremal charged black hole $\sim$ D-branes
- Excitations can be counted, account for

$$S = \frac{A}{4G_N}$$

[Strominger-Vafa], [Maldacena-Strominger-Witten]

- Extremal charged black hole $\sim$ D-branes
- Excitations can be counted, account for

$$S = \frac{A}{4G_N} + \dots$$

[Strominger-Vafa], [Maldacena-Strominger-Witten]

- Also predicts subleading corrections.

Corrections to the area law

- Einstein-Hilbert

$$\mathcal{L} = \sqrt{-g} \frac{R}{16\pi G_N}$$

- Area law

$$S = \frac{A}{4G_N}$$

Corrections to the area law

- Einstein-Hilbert corrected:

$$\mathcal{L} = \sqrt{-g} \left(\frac{R}{16\pi G_N} + \frac{c}{2} R^2 + \dots \right)$$

- Area law accordingly modified :

$$S = \frac{A}{4G_N} + 8\pi c \int_{\text{hor}} R_{rtrt} + \dots$$

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$$S = \frac{A}{4G_N} + 8\pi c \int_{\text{hor}} R_{rtrt} + \dots$$

- Wald's formula:

$$S = -2\pi \int_{\text{hor}} \sqrt{-g} \frac{\delta \mathcal{L}}{\delta R_{abcd}} \epsilon_{ab} \epsilon_{cd}$$

String theory works

- Two way to calculate corrections:

Microscopic

d.o.f. on the brane

Macroscopic

Wald's formula

- Completely agrees !

[de Wit et al.][Ooguri-Strominger-Vafa]

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[de Wit et al.][Ooguri-Strominger-Vafa]

- In four dimensions.

Odd dimensions

- Black rings in 5d.
- Entropy correction:

Microscopic

done

Macroscopic

not yet

- Why ?

Odd dimensions

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- Why ? Wald's formula **isn't** applicable to gravitational Chern-Simons:

$$\int F \wedge \text{tr } \Gamma \wedge R$$

Odd dimensions

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- Entropy correction:

Microscopic

done

Macroscopic

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Our aim today

Entropy correction from grav. CS.

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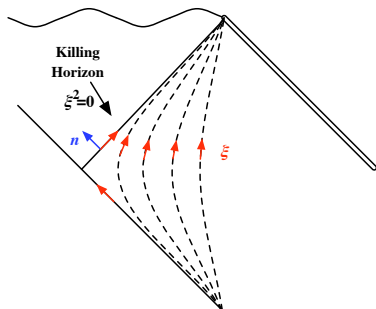
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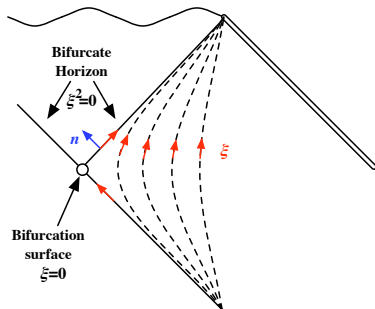
- Killing Horizon : $\xi^2 = 0$.
- Binormal

$$\epsilon_{ab} = \xi_a n_b - \xi_b n_a$$

- Surface gravity

$$\nabla_a \xi_b = \kappa \epsilon_{ab} + \dots$$

Terminology

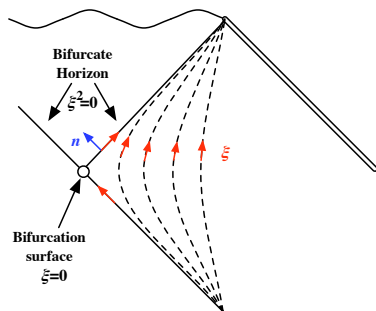


- Bifurcation Surface

$$\xi = 0.$$

- Bifurcate horizon :
A pair of horizons
which pass BS.

Hawking radiation

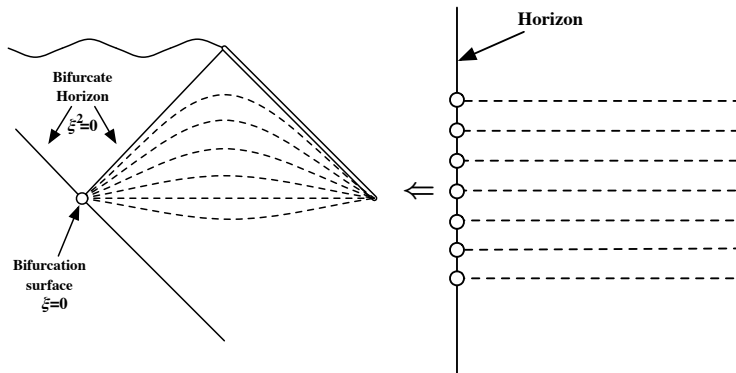


- Free field radiates at

$$T_H = \kappa/2\pi.$$

- Unruh effect near B.S.
- Depends only on bg metric,
- Not quantum gravity per se.

Bifurcation surface



- Horizon cross section at finite t **is** the b.s.

- At the B.S.,

$$\nabla_c \epsilon_{ab} = \kappa^{-1} \nabla_c \nabla_a \xi_b = \kappa^{-1} R_{abcd} \xi^d = 0.$$

→ holonomy reduces to

$$SO(D-1, 1) \supset SO(1, 1)_N \times SO(D-2)$$

Black Hole Thermodynamics

Zeroth law

κ constant on the horizon.

First law

$$\frac{\kappa}{2\pi} \frac{\delta A}{4G_N} = \delta m - \Omega \delta J$$

Second law

A increases with time.

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Higher derivative corrections

- RG perspective
 - Planck suppressed terms R^2 , R^4 etc.
- Coefficients calculable in string compactifications
- Affects black hole solutions.

Conceptual problems

- Null dir. of $g_{ab} \neq$ maximal propagation speed.
- e.g. Field redefinition

$$g_{ab} \rightarrow g_{ab} + cR_{ab} + c'Rg_{ab} + \dots$$

changes the 'light cone'

- Physics should be invariant !

Bifurcate Horizon in Higher derivative gravity

- Define the horizon using the light cone of g_{ab} .
- Suppose it has Killing, bifurcate horizon :
under $g_{ab} \rightarrow g_{ab} + h_{ab}$,

Bifurcation surface

invariant, defined by $\xi^a = 0$

Bifurcate horizon

$$\xi^a \xi_a \rightarrow g_{ab} \xi^a \xi^b + h_{ab} \xi^a \xi^b.$$

2nd term invariant under ξ

→ zero by evaluating at B.S.

Hawking temperature

Evaluate $\kappa^2 = |\nabla_a \xi^b|^2$ at B.S.

→ Christoffel drops off because $\xi^a = 0$.

Black Hole Thermo. in Higher Derivative Gravity

Zeroth law

κ constant on the horizon : assumption

First law

$$\frac{\kappa}{2\pi} \delta S = \delta m - \Omega \delta J, \quad S \text{ given by Wald's formula}$$

Second law

Difficult to establish

Black Hole Entropy

$$S = -2\pi \int_{\text{hor}} \sqrt{-g} \frac{\delta \mathcal{L}}{\delta R_{abcd}} \epsilon_{ab} \epsilon_{cd}$$

for $\mathcal{L} = \mathcal{L}(g_{ab}, R_{abcd}, \nabla_e R_{abcd}, \dots; \phi, \dots)$

- It does not include gravitational Chern-Simons

$$*\delta \mathcal{L} = \text{tr } \Gamma \wedge R^{2n-1}$$

- or Green-Schwarz type coupling

$$\delta \mathcal{L} = \frac{1}{2} |H|^2 \quad \text{with} \quad H = dB + \text{tr } \Gamma \wedge R.$$

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D space-time dimension

$L(\phi)$ Lagrangian density as D -form

ϕ collective symbol for the fields. $g_{\mu\nu}, A_\mu, \dots$

$\mathcal{L}_\xi$ Lie derivative by a vector field ξ ,

$$\mathcal{L}_\xi \omega = (d\iota_\xi + \iota_\xi d)\omega.$$

δ_ξ Variation under diffeo. e.g.

$$\delta_\xi \Gamma_b^a = \mathcal{L}_\xi \Gamma_b^a + d(U_\xi)_b^a$$

where $(U_\xi)_b^a = \partial_b \xi^a$.

Diffeomorphism Invariance

- Assumption by Wald:

$$\delta_\xi L(\phi) = \mathcal{L}_\xi L(\phi).$$

Cannot incorporate the CS.

- Today:

$$\delta_\xi L(\phi) = \mathcal{L}_\xi L(\phi) + d\Xi_\xi.$$

- Almost verbatim transcript of Wald's original.

Covariant Hamiltonian Method

- EOM E_ϕ and symplectic potential Θ via

$$\delta L = E_\phi \delta \phi + d\Theta(\phi, \delta \phi)$$

- Θ pairs of ‘coordinate’ and ‘momenta’

$$L(\phi) = \frac{1}{2} * d\phi \wedge d\phi \rightarrow \Theta = *d\phi \wedge \delta \phi$$

- symplectic form given by

$$\Omega(\phi, \delta_1 \phi, \delta_2 \phi) = \delta_1 \Theta(\phi, \delta_2 \phi) - \delta_2 \Theta(\phi, \delta_1 \phi).$$

Lemma

- j : a form constructed from
 - fields ϕ
 - an external field ξ
- suppose j is closed on-shell for any ξ .
- It is then exact on-shell. i.e.

$$dj_\xi \simeq 0 \rightarrow j_\xi \simeq dQ_\xi.$$

- $\simeq$: equality on-shell.

Noether's theorem

- Symmetry leads to conserved current :

$$j_\xi = \Theta(\phi, \delta_\xi \phi) - \iota_\xi L - \Xi_\xi$$

satisfies

$$dj_\xi \simeq 0$$

- The lemma implies

$$j_\xi \simeq dQ_\xi, \quad \rightarrow \quad \int_C j_\xi \simeq \int_{\partial C} Q_\xi.$$

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i.e.

Global symmetry	$\rightarrow$	Conserved charges
Gauge symmetry	$\rightarrow$	Gauss law

Un-illuminating main part

- Define Π_ξ via

$$\delta_\xi \Theta = \mathcal{L}_\xi \Theta + \Pi_\xi.$$

- Recall $\delta L = E_\phi \delta \phi + d\Theta$, $\delta_\xi = \mathcal{L}_\xi L + d\Xi_\xi$.

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- Define Π_ξ via

$$\delta_\xi \Theta = \mathcal{L}_\xi \Theta + \Pi_\xi.$$

- Recall $\delta L = E_\phi \delta \phi + d\Theta$, $\delta_\xi = \mathcal{L}_\xi L + d\Xi_\xi$.
- Calculating $\delta \delta_\xi L$ in two ways:

$$d\Pi_\xi \simeq \delta d\Xi_\xi \longrightarrow \Pi_\xi - \delta \Xi_\xi \simeq d\Sigma_\xi.$$

$$\begin{aligned} \longrightarrow \delta j_\xi &= \delta \Theta(\phi, \delta_\xi \phi) - \iota_\xi \delta L - \delta \Xi_\xi \\ &\simeq \delta \Theta(\phi, \delta_\xi \phi) - \delta_\xi \Theta(\phi, \delta \phi) + d\iota_\xi \Theta + \Pi_\xi - \delta \Xi_\xi \\ &\simeq \Omega(\phi, \delta \phi, \delta_\xi \phi) + d(\iota_\xi \Theta + \Sigma_\xi). \end{aligned}$$

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- Thus, for C_ξ with $\delta C_\xi = \iota_\xi \Theta + \Sigma_\xi$,

$$\delta dQ'_\xi \simeq \Omega(\phi, \delta \phi, \delta_\xi \phi)$$

where

$$Q'_\xi = Q_\xi - C_\xi.$$

Recap.

- $\delta dQ'_\xi \simeq \Omega(\phi, \delta\phi, \delta_\xi\phi)$ means

$\int dQ'_\xi$ is the Hamiltonian generating ξ .

- Let $\xi = t + \Omega\phi$ where
 - ξ Horizon generating Killing
 - t global time translation
 - ϕ angular rotation

Then

$$\delta \int_{\text{hor}} Q'_\xi \simeq \delta \int_{\infty} Q'_t + \Omega \delta \int_{\infty} Q'_\phi.$$

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- $E = \int_{\infty} Q'_t$, $J = \delta \int_{\infty} Q'_\phi$.
- Taking the horizon at the bifurcation surface,

$$\int_{\text{hor}} Q'_\xi = \kappa \int_{\text{hor}} Q'_\xi \big|_{\xi \rightarrow 0, \nabla_a \xi_b \rightarrow \epsilon_{ab}}$$

- Then, define

$$S = 2\pi \int_{\text{hor}} Q'_\xi \big|_{\xi \rightarrow 0, \nabla_a \xi_b \rightarrow \epsilon_{ab}}$$

Result.

$$\frac{\kappa}{2\pi} \delta S = \delta E + \Omega \delta J.$$

3d Chern-Simons

- Start from $L_{CS} = \beta \operatorname{tr}(\Gamma R - \frac{1}{3}\Gamma^3)$
 $\rightarrow \delta_\xi L_{CS} = \mathcal{L}_\xi L_{CS} - d(\beta \operatorname{tr} dU_\xi \Gamma) \rightarrow \Xi_\xi = -\beta \operatorname{tr} dU_\xi \Gamma$
- Non-covariant part in Θ is $-\beta \operatorname{tr} \Gamma \delta \Gamma$
 $\rightarrow \Pi_\xi = -\beta \operatorname{tr} dU_\xi \delta \Gamma$
 $\rightarrow \Pi_\xi - \delta \Xi_\xi \simeq 0 \rightarrow \Sigma_\xi = 0$
 $\rightarrow j_\xi = \Theta(\phi, \delta_\xi \phi) - \Xi_\xi + \dots = 2\beta \operatorname{tr} dU_\xi \Gamma + \dots$
 $\rightarrow Q'_\xi = 2\beta \operatorname{tr} U_\xi \Gamma + \dots$
- $(U_\xi)_b^a = \partial_b \xi^a$.

3d Chern-Simons

- $L_{CS} = \beta \operatorname{tr}(\Gamma R - \frac{1}{3}\Gamma^3) \rightarrow S_{CS} = 8\pi\beta \int_{\text{hor}} \Gamma_N$

where $\Gamma_N = -\epsilon^\nu{}_\mu \Gamma^\mu{}_{\nu\rho} dx^\rho / 2$.

- agrees with known corrections to BTZ.
- Naive application of Wald $\rightarrow 4\pi\beta \int_{\text{hor}} \Gamma_N$.

Chern-Simons in General Dimensions

- $L_{CS} = \beta \operatorname{tr}(\Gamma R^{2m-1} + \dots)$

→ $S_{CS} = 8\pi m\beta \int_{\text{hor}} \Gamma_N R_N^{2m-2}$

- **Recall**

Holonomy at the bifurcation surface

$$= SO(1, 1)_N \times SO(D - 2)$$

- **Entropy correction = CS of the normal bundle !**

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Summary

- **Spacetime with black holes reviewed.**
- **Entropy correction from the grav. CS.**

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- Spacetime with black holes reviewed.
- Entropy correction from the grav. CS.

Outlook – Need application !

- Black rings. – working on it with friends.
- 7d black holes. – idea wanted.

$$\text{tr } \Gamma \wedge R^3 \longrightarrow \int_{\text{hor}} \Gamma_N \wedge R_N^2$$